

Fall 2014

## Optics

Part I:

Round trip phase delay:  $\delta_1 = \frac{4\pi\nu_1 l}{c} + 2\phi_{r1} = m\pi$

$$\delta_2 = \frac{4\pi\nu_2 l}{c} + 2\phi_{r2} = m\pi$$

Equating  $\delta_1$  &  $\delta_2$ :  $\frac{4\pi(\nu_1 - \nu_2)l}{c} + 2(\phi_{r1} - \phi_{r2}) = 0$

$$(\phi_{r1} - \phi_{r2}) = \frac{2\pi l}{c} (\nu_1 - \nu_2) =$$

$$\frac{2\pi (1.5)}{3 \times 10^8} (10^7) = 2\pi (0.05)$$

$$\therefore a\lambda = 0.05\lambda = \lambda/20$$

Alternatively: A  $\frac{\lambda}{2}$  plate would produce coincident

mode frequencies  $l = m\frac{\lambda_1}{2} = \frac{(m+1)\lambda_2}{2}$

$$\therefore \frac{\lambda}{2} \text{ retardation produces } \Delta\nu = \frac{c}{2l} = 100 \text{ MHz.}$$

For  $\Delta\nu = 10 \text{ MHz}$ , we must have

$$\boxed{a\lambda = \left(\frac{10}{100}\right) \left(\frac{\lambda}{2}\right) = \frac{\lambda}{20}}$$

## Part 2:

At low power,  $I(\nu_0) = I_0 \exp[-\alpha(\nu_0)l]$

$$\Rightarrow \alpha(\nu_0)l = 0.2 \quad (\text{Power absorption})$$

But  $E = E_0 \exp[-\alpha(\nu_0) + j\beta(\nu_0)]l$  [amplitude]

where  $\alpha(\nu_0) = \frac{\alpha(\nu_0)}{2} \leftarrow \text{Intensity}$   
Ampl.  $\therefore \alpha(\nu_0)l = 0.1$

We have  $\frac{kX''(\nu_0)l}{2n^2} = -\alpha(\nu_0)l = 0.1$

From Kramers-Kronig relations,  $X' = \frac{X''_2(\nu_0 - \nu)}{\Delta\nu}$

$$\Rightarrow X'(\nu_c) = X''(\nu_c) \frac{2(\nu_0 - \nu_c)}{2(\nu_0 - \nu_c)} = X''(\nu_c)$$

Def of  $\Delta\nu$

Also,  $X''(\nu_c) = \frac{X''(\nu_0)}{2}$  (by def. of  $\nu_c$ )

$$\therefore \beta(\nu_c)l = \frac{kX'(\nu_c)l}{2n^2} = \frac{kX''(\nu_c)l}{2n^2} = \frac{kX''(\nu_0)l}{4n^2} = \frac{\alpha(\nu_0)l}{2} = 0.05$$

This corresponds to a single pass delay. For the round-trip delay,  $\delta = 2(0.05) = 0.1$ . This is the phase shift at low power. When the transition saturates, this phase shift goes to zero, and detunes the Fabry-Perot

$$T = \frac{(1-R)^2}{(1-R)^2 + 4R \sin^2(\delta/2)} = \frac{(1-0.9)^2}{(1-0.9)^2 + 4(0.9)(0.05)^2} = 0.52$$